

Phase relations of Fe-Ni alloys at high pressure and temperature

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Abstract

Using a diamond anvil cell and double-sided laser-heating coupled with synchrotron x-ray diffraction, we determined phase relations for three compositions of Fe-rich FeNi alloys *in situ* at high pressure and high temperature. We studied Fe with 5 wt%, 15 wt%, and 20 wt% Ni to 55 GPa, 62 GPa, and 72 GPa respectively at temperatures up to ~3000 K. Ni stabilizes the face centered cubic phase to lower temperatures and higher pressure, and this effect increases with increasing pressure. Extrapolation of our experimental results for Fe with 15 wt% Ni suggests that the stable phase at inner core conditions is hexagonal close packed, although if the temperature at the inner core boundary is higher than ~6400 K, a two phase outer region may also exist. Comparison to previous laser-heated diamond anvil cell studies demonstrates the importance of kinetics even at high temperatures.

Keywords: Iron-Nickel alloys; Earth's core; High pressure; Diamond anvil cell; Laser-heating; X-ray Diffraction

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1 **1. Introduction**

2 As major constituents of the Earth's core, Fe and its alloys have long been of great
3 interest to geophysicists (Birch, 1952). From geophysical, cosmochemical, and
4 geodynamic data, it is generally accepted that the core is composed of an Fe-rich Fe-Ni
5 alloy with a small amount of light element(s) (McDonough, 2003; Stixrude and Brown,
6 1998). Information on the behavior of Fe-Ni alloys at high pressure-temperature (P - T),
7 such as phase relations and thermal equations of state (EOS), is essential for interpreting
8 seismic and geomagnetic observations and for numerical modeling of the Earth's deep
9 interior.

10 While pure Fe at high P - T has been the subject of numerous studies, there are far
11 fewer studies on Fe-Ni alloys. At ambient conditions, the stable phase of Fe is the body-
12 centered cubic (bcc) structure. This phase, α -Fe, transforms into a face-centered cubic (fcc
13 or γ -Fe) phase upon increasing temperature above 1185 K, and then transforms to another
14 bcc (δ -Fe) phase before melting (Hemley and Mao, 2001). At high pressure, both bcc and
15 fcc Fe transform into the hexagonal close-packed (hcp or ϵ) phase (Takahashi and Bassett,
16 1964). This phase has a broad stability region, and Fe is generally considered to be stable
17 in the hcp structure at inner core conditions (Hemley and Mao, 2001; Ma et al., 2004).

18 Previous work on the quasihydrostatic compression of $\text{Fe}_{0.8}\text{Ni}_{0.2}$ (Mao et al., 1990)
19 are the only results to core pressures, but they were measured at ambient temperature. At
20 high pressure Ni has been shown to stabilize the face-centered cubic (fcc) phase (Huang et
21 al., 1988; Lin et al., 2002). Externally heated diamond anvil cell (DAC) studies on a
22 number of Fe-Ni alloys with up to 35 wt% Ni indicate that the core may consist of both a
23 fcc and hcp phase (Huang et al., 1988; Huang et al., 1992), but these studies require

24 significant extrapolation for application to the core. Recent laser-heated DAC (LHDAC)
25 results on Fe with 10 wt% Ni were conducted to much higher temperatures and suggest
26 that the inner core should be in the hcp structure (Lin, 2002). It is necessary to further
27 explore the compositional effect of Ni on the P - T - x phase diagram of Fe-Ni alloys at high
28 P - T and relevant core compositions to determine the effect on core properties. This would
29 serve as a baseline for understanding the additional effects of the presence of potential light
30 elements in the core.

31

32 **2. Experimental**

33 Three compositions of Fe-rich Fe-Ni alloys were investigated. Fe15%Ni and
34 Fe20%Ni refer to synthetic alloys of Fe with 15 wt% (14.4 mol%) Ni and 20 wt% (19.2
35 mol%) Ni respectively. In this paper Fe5%Ni refers to filings from a natural sample of the
36 Negrillos (IIAB) iron meteorite which contains 5.4 wt% (5.2 mol%) Ni. This specimen
37 also contains minor amounts of other elements, most notably Co (0.4 wt%) and P (0.1-0.2
38 wt%). This composition is probably similar to the Earth's core minus the light element
39 component (McDonough, 2003). Small foils of sample (7 μm thick, 50 μm diameter) were
40 sandwiched between thin (7 μm) layers of NaCl which acted as a thermal insulator,
41 pressure-transmitting medium, and internal pressure standard (Heinz and Jeanloz, 1994;
42 Sata et al., 2002). The pressure was calculated from the ambient temperature NaCl EOS
43 before heating, similar to the previous Fe-Ni alloy LHDAC study (Lin et al., 2003). The
44 sample assemblage was loaded into a \sim 100 μm diameter chamber that was drilled into a
45 pre-indented stainless steel gasket in a symmetric DAC.

46 Monochromatic synchrotron x-radiation ($\lambda = 0.3344 \text{ \AA}$) at GSECARS 13-IDD
47 beamline of the Advanced Photon Source (APS), Argonne National Laboratory (ANL) was
48 used for angle-dispersive x-ray diffraction measurements. The x-ray beam was focused in
49 the horizontal and vertical direction to a $5 \times 7 \text{ \mu m}$ (FWHM) spot using Kirkpatrick-Baez
50 mirrors. Diffraction patterns were taken using a Bruker CCD detector *in situ* at high P - T . A
51 CeO_2 standard was used to calibrate the sample to detector distance and the detector tilt.
52 An x-ray transparent cubic boron nitride seat was used to support the downstream diamond
53 in order to increase 2θ access from the sample down to d-spacings of 0.8 \AA .

54 A double-sided Nd:YLF laser system operating in the Gaussian TEM_{00} mode was
55 used to heat the sample (Shen et al., 2001). Temperature was determined using spectral
56 radiometry by fitting the thermal radiation spectrum to the Planck radiation function.
57 Temperature measurements taken by the spectral radiometry system at 13-IDD have been
58 verified prior to our experiments. The major source of potential error in temperature is the
59 lack of information regarding the wavelength dependence of emissivity. The wavelength
60 range used for the temperature fit was 670-830 nm which represents a compromise
61 between reducing the error from emissivity by reducing the wavelength range, while also
62 having a large enough range to avoid introduction of errors from some non-linear portion
63 of the system response (Shen et al., 2001). The laser beam diameter was approximately 25
64 μm , producing a stable heated spot much larger than the focused x-ray beam. Axial and
65 radial temperature gradients were minimized by the use of the NaCl insulating layers. The
66 maximized heating area and minimized temperature gradients help reduce errors associated
67 with chromatic aberration. The temperature uncertainty (1σ) for multiple temperature

68 measurements and the temperature difference measured from both sides of the sample was
69 approximately 100 K.

70

71 **3. Results and Discussion**

72 We investigated the phase boundaries for the hcp + fcc two phase coexistence
73 region by monitoring changes in the x-ray diffraction pattern *in situ* as we laser-heated the
74 sample at high P . Each heating cycle was conducted on a new, previously unheated spot to
75 ensure that the composition of the spot probed was the same as the composition of the
76 sample. Upon increasing temperature, we observed the transformation from hcp only to a
77 mixture of hcp + fcc phases to fcc (Figure 1). The hcp \rightarrow hcp + fcc and hcp + fcc \rightarrow fcc
78 phase boundaries for varying Ni compositions were determined for the three compositions
79 studied (Figure 2). In the two phase region, the volume difference between the Ni-rich fcc
80 phase and the coexisting Ni-poor hcp phase ($\Delta V^{\text{fcc-hcp}}$) was approximately 1-3% with the
81 difference becoming larger at higher P - T (Table 1). This corresponds to an approximate
82 10 wt% difference in Ni content between the fcc and hcp phases (Figure 3). A previous
83 study on Fe10%Ni found a 1% $\Delta V^{\text{fcc-hcp}}$ at 40 GPa and 1602 K (Lin et al., 2002).

84 Reproducibility of results depends on reaching thermodynamic equilibrium. We
85 investigated the importance of kinetics during a number of long heating cycles. For
86 example, it took nearly 30 minutes of laser-heating at ~ 1400 K to transform the Fe20%Ni
87 sample that was initially a quenched mixture of hcp + fcc phases into hcp only (Figure 5).
88 Even at high temperature, kinetic effects indicate that a significant time may be required
89 for the sample to reach its equilibrium phase assemblage. This may explain why our hcp +
90 fcc coexistence region is narrower than those of the previous laser-heating study (Lin et al.,

91 2002). Diffusion of Ni appears to be an important issue. Upon decreasing temperature, it
92 was very difficult to completely transform the quenched fcc phase back into the hcp phase.
93 This stubborn nature of the fcc phase has been observed previously in pure Fe and other
94 Fe-rich alloys (Uchida et al., 2001). We mainly determined transition conditions on
95 increasing temperature cycles, but did constrain the phase boundaries on some long heating
96 cycles following a decrease in temperature (Figure 6).

97 Increasing Ni concentrations stabilize the fcc phase to lower T and higher P . Even 5
98 wt% Ni shows significant stabilization of the fcc field relative to pure Fe, and this effect
99 increases with increasing pressure. Figure 3 shows a summary P - T - x phase diagram for Fe-
100 rich Fe-Ni compositions. Qualitatively, our results are consistent with previously reported
101 work on different Fe-Ni compositions (Huang et al., 1988; Huang et al., 1992; Lin et al.,
102 2002) and the results for pure Fe (Ma et al., 2004). However, earlier resistively heated
103 DAC studies at higher Ni (up to 35 wt%) concentrations report the hcp + fcc \rightarrow fcc phase
104 boundary at much lower temperature than would be expected from extrapolation of our
105 data to higher Ni content (Huang et al., 1992). Additional compositions may need to be
106 measured in order to constrain the Fe-Ni phase diagram with higher Ni contents. While a
107 resistively heated DAC provides better temperature control, it is limited to lower
108 temperatures, so the hcp + fcc \rightarrow fcc boundary has not been reported using this technique
109 for lower Ni concentrations. We observed that the slope of the hcp \rightarrow hcp + fcc and hcp +
110 fcc \rightarrow fcc phase boundaries becomes flatter with increasing Ni content which is consistent
111 with previous results, although extrapolation between our 15 wt% Ni data predict a
112 shallower slope than reported for Fe10%Ni (Lin et al., 2002) (Table 2).

113 Constraining the phase boundary slopes is critical when considering extrapolation
114 of the experimental results to core conditions. The Lindemann law extrapolation of recent
115 LHDAC data gives a melting temperature for pure Fe of 5800 K at the pressure of the
116 inner core boundary (ICB) (Ma et al., 2004). Extrapolation of Fe10%Ni and Fe15%Ni
117 results indicate that hcp would be the stable phase for these compositions for possible core
118 geotherms crossing through 5800 K at ICB pressure (Figure 4). It should be noted that
119 there are many uncertainties in our understanding of the actual ICB temperature. For
120 example, the addition of light elements would depress the melting point and result in a
121 lower ICB temperature (Alfé et al., 2002). However, if our simple linear extrapolation
122 holds, the ICB would have to be ~6400 K for the two phase hcp + fcc to be stable for
123 Fe15%Ni at ICB pressure. If the ICB were at even higher temperature this would stabilize
124 the two phase hcp + fcc region for even lower Ni contents. Then the inner core could
125 comprise an outer two phase hcp + fcc layer and a hcp only innermost core. Whether this
126 has relevance to the observed change in elastic anisotropy in the inner core (Ishii and
127 Dzierwonski, 2002) requires further investigation and tighter constraints of the ICB
128 temperature.

129 Theoretical calculations of pure Fe have resulted in interesting speculations
130 regarding the inner core. Two recent studies suggested that the bcc structure of Fe may be
131 the stable inner core phase (Belonoshko et al., 2003; Vocadlo et al., 2003). It is known from
132 the 1 bar phase diagram that alloying with Ni promotes the fcc phase and suppresses the bcc
133 phase in Fe-rich alloys, and as shown in this study, Ni stabilizes the fcc phase even to high
134 *P-T* conditions. We expect that the presence of Ni would counteract the formation of the

135 hypothesized bcc phase at high P - T , but the effect of Ni on the stability of phases should be
136 investigated experimentally to inner core conditions to confirm this.

137

138 **4. Conclusions**

139 The addition of Ni clearly affects the physics and chemistry of Fe at high P - T . In
140 examining phase relations in Fe5%Ni, Fe15%Ni, and Fe25%Ni we found that increasing
141 Ni content stabilizes the fcc phase to lower temperature and higher pressures. Our results
142 are consistent with previous LHDAC studies, but discrepancies with resistively heated
143 DAC studies remain. This may be an issue with kinetics, which was found to be important
144 especially with the fcc phase. Fe-rich Fe-Ni alloys should serve as a baseline composition
145 for comparison to seismic observations and the study of possible light alloying elements.
146 Ni may influence partitioning of minor and trace elements between the solid, inner core
147 and the liquid, outer core. Its effect on the melting curve of Fe also needs to be studied. In
148 general, extension of LHDAC studies to higher P - T is needed. In particular, key anchor
149 points for extrapolation of the melting curve and determination the phase(s) in the inner
150 core would be the intersection of the hcp \rightarrow hcp + fcc boundary and the melting curve and
151 the intersection of the hcp + fcc \rightarrow fcc boundary and the melting curve.

152

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Figure Captions

Figure 1. Representative XRD spectra for Fe5%Ni at 43 GPa. The sample transforms from an initially hcp phase to a mixture of hcp and fcc phases and finally to nearly entirely fcc phase with increasing temperature. Diffraction peaks from the different phases as well as those from the B2 phase of NaCl are labeled.

Figure 2. Phase relationships for Fe5%Ni, Fe15%Ni, and Fe20%Ni. The solid lines show approximate phase boundaries for the hcp + fcc two phase coexistence region. Symbols show experimental results where we observed different phases; *:hcp only; squares:hcp+fcc; triangles:fcc only; X:melt. Lines may differ from thermodynamic equilibrium due to kinetics.

Figure 3. Schematic phase relationships for Fe-rich Fe-Ni compositions for varying $T-x$ at 30, 40, 50, and 60 GPa. Circles: interpolations of our results for Fe5%Ni, Fe15%Ni, and Fe20%Ni; *: Fe (Ma et al., 2004); triangles:Fe10%Ni (Lin et al., 2002); diamonds:Fe30%Ni (Huang et al., 1992). Closed symbols and dotted lines show hcp→hcp + fcc boundary; open symbols and solid lines show hcp + fcc→fcc boundary.

Figure 4. Portion of XRD spectra for Fe20%Ni taken at 72 GPa and 1400 K at different times during a 50 minute heating cycle. At right, Bruker CCD images show that the Debye rings became much smoother with extended heating, indicating the decreasing grain size of the hcp phase equilibrating at high temperature relative to the grain size in the hcp + fcc

assemblage initially quenched from higher temperature. It took approximately 30 minutes to transform the initial mixed hcp + fcc assemblage into hcp only.

Figure 5. Conditions for the long heating cycle for Fe20%Ni at 72 GPa corresponding to Figure 4. Before this cycle, a hcp+fcc two phase mixture was synthesized at 1596 K. Dotted line indicates when power to YLF lasers was increased. Symbols show what phases were present in the *in-situ* XRD patterns taken at different times during heating. We crossed the hcp only \rightarrow hcp + fcc phase boundary on both decreasing and increasing temperature to constrain this boundary within 100 K. Dashed line indicates approximate transition temperature.

Figure 6. Extrapolation of our Fe15%Ni and previously reported Fe10%Ni (Lin et al., 2002) results to core conditions. Thick solid line: Lindemann law extrapolation of hcp Fe melting curve (Ma et al., 2004); Dotted lines: hcp \rightarrow hcp + fcc boundaries; thin, solid lines: hcp + fcc \rightarrow fcc boundaries Shaded areas indicate two phase hcp + fcc regions. Dashed lines show the core-mantle boundary (CMB) and inner core-outer core boundary (ICB) pressures.

Table 1. $\Delta V^{\text{fcc-hcp}}$ for different compositions and P - T conditions

Composition	P (GPa)	T (K)	$\Delta V^{\text{fcc-hcp}}$ (%)
Fe5%Ni	43	1680	1.3
	55	2280	2.4
Fe15%Ni	39	1450	1.1
	53	1740	3.2
Fe20%Ni	62	1786	2.2
	72	1909	2.2

Table 2. Phase boundary slope for varying Ni compositions

Slope (K/GPa)	Pure Fe	Fe10%Ni	Fe15%Ni	Fe30%Ni
hcp \rightarrow fcc	39.3	-	-	-
hcp \rightarrow hcp + fcc	-	28.1	17.9	-
hcp + fcc \rightarrow fcc	-	34.9	21.8	25

Fe10%Ni data is from (Lin et al., 2002); Fe30%Ni data from (Huang et al., 1992); Note: there was not enough data to estimate a phase boundary for Fe5%Ni and Fe20%Ni.

Figure 1
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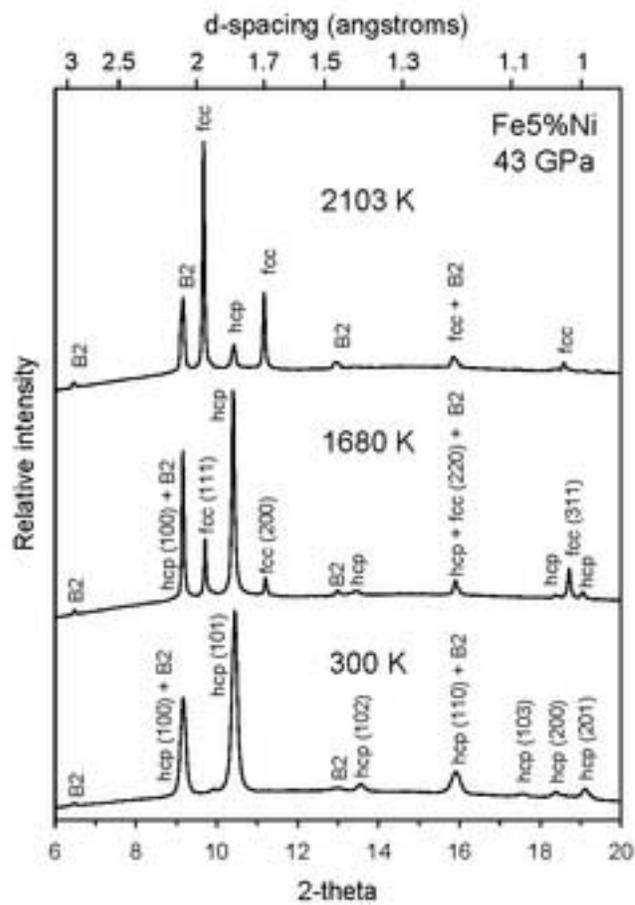


Figure 1

Figure 2

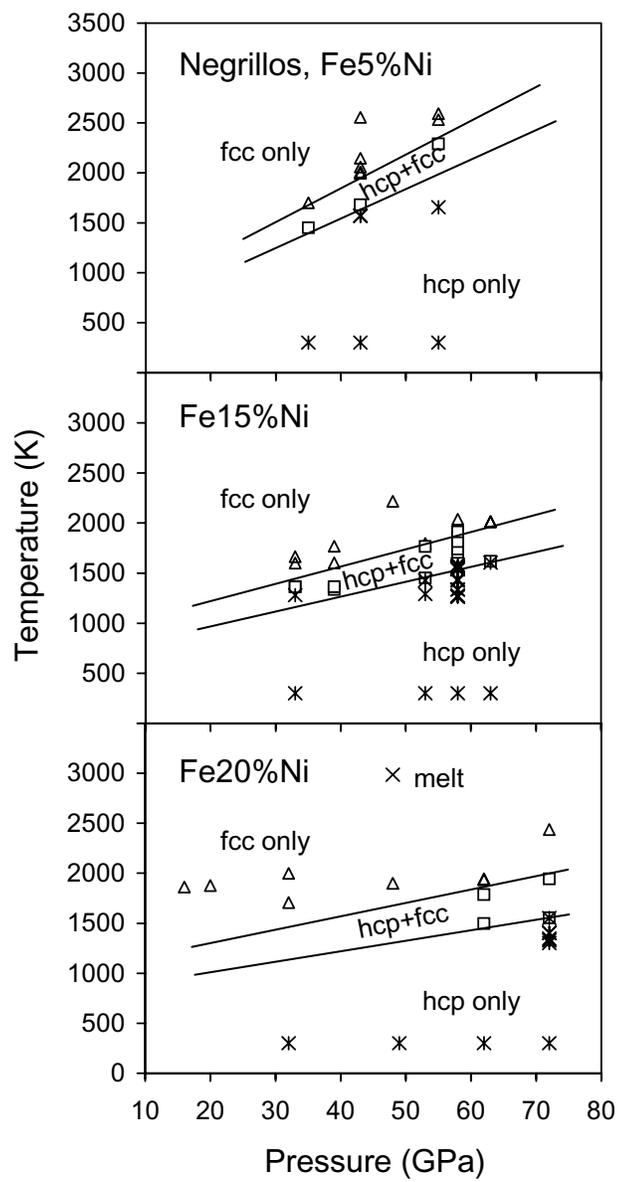


Figure 2

Figure 3

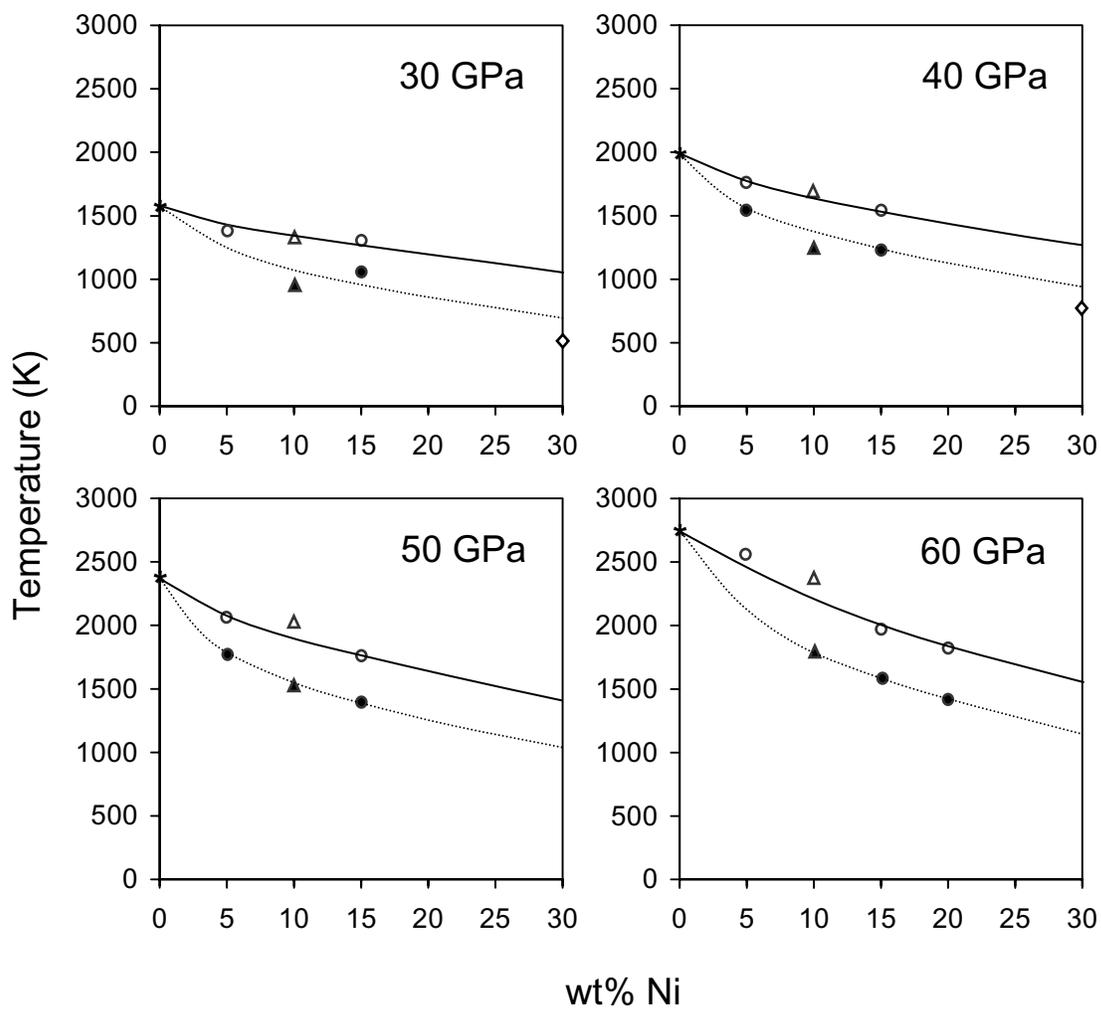


Figure 3

Figure 4
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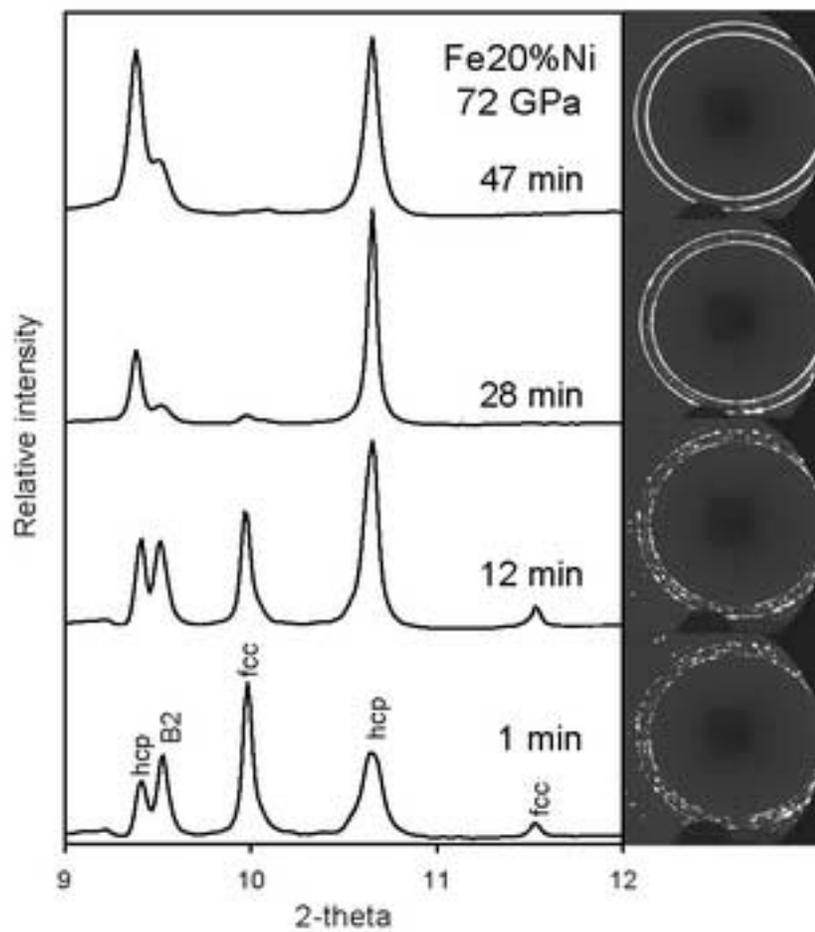


Figure 4

Figure 5

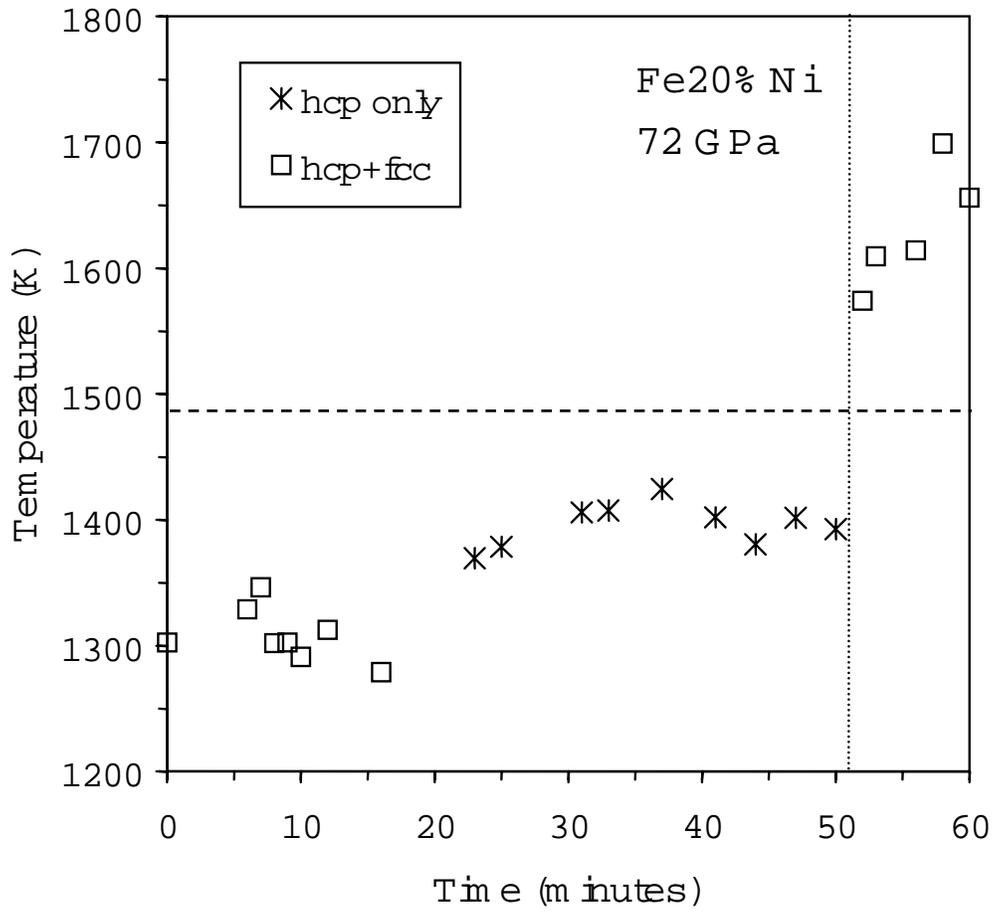


Figure 5

Figure 6
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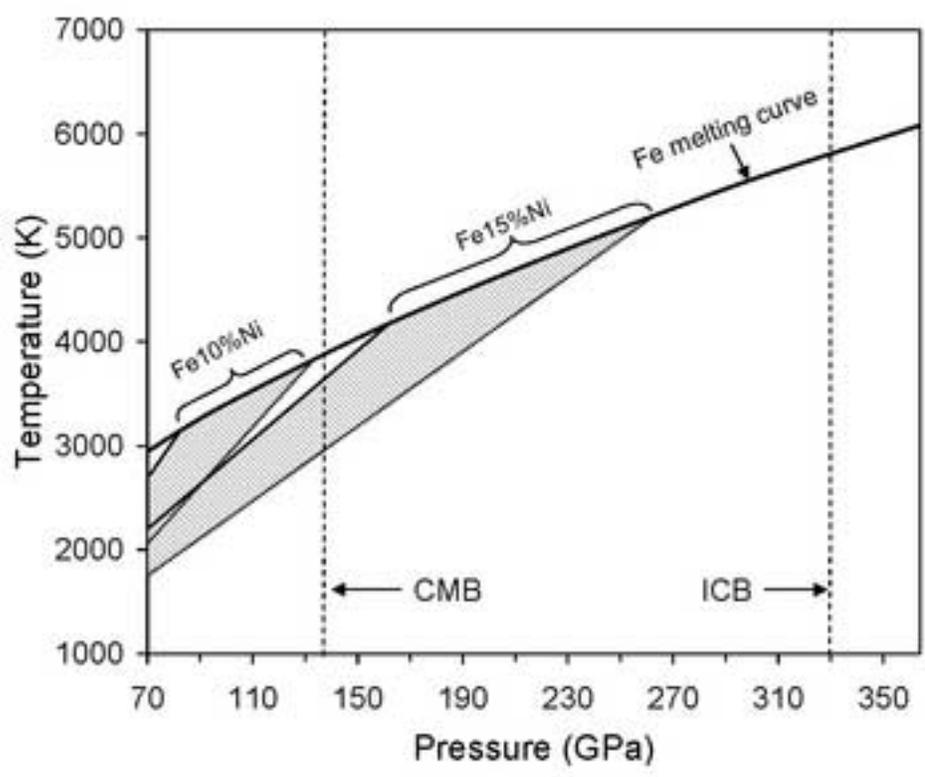


Figure 6